



Chap 6: Manipulator Dynamics

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Acceleration of a Rigid Body -1

- Differentiation of a velocity vector V_Q

$${}^B A_Q = \frac{d}{dt} {}^B V_Q = \lim_{\Delta t \rightarrow 0} \frac{{}^B V_Q(t + \Delta t) - {}^B V_Q(t)}{\Delta t}$$

Derivative of velocity vector ${}^B V_Q$ relative to frame $\{B\}$

$${}^A ({}^B A_Q) = {}^A \left(\frac{d}{dt} {}^B V_Q \right)$$

Expressed in frame $\{A\}$

$$= \underline{{}^A R} {}^B ({}^B A_Q) = \underline{{}^A R} {}^B A_Q$$

When both frames are the same

$$a_C = {}^U A_{C ORG}$$

Acceleration of the origin of frame $\{C\}$ relative to the universe reference frame $\{U\}$

Acceleration of a Rigid Body -2

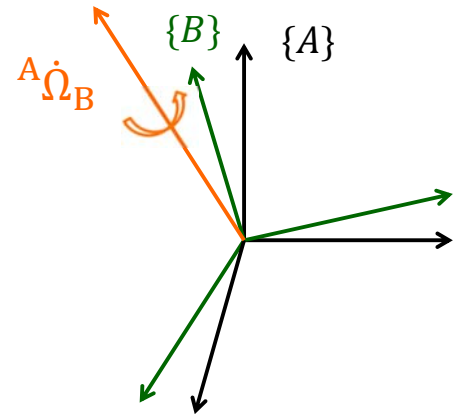
- Angular acceleration vector ${}^A\dot{\Omega}_B$

$${}^A\dot{\Omega}_B = \frac{d}{dt} {}^A\Omega_B = \lim_{\Delta t \rightarrow 0} \frac{{}^A\Omega_B(t + \Delta t) - {}^A\Omega_B(t)}{\Delta t}$$

Derivative of angular velocity of frame $\{B\}$ relative to frame $\{A\}$

$${}^C({}^A\dot{\Omega}_B)$$

Expressed in frame $\{C\}$



$$\dot{\omega}_C = {}^U\dot{\Omega}_C$$

Angular acceleration of frame $\{C\}$ relative to the universe reference frame $\{U\}$

Acceleration of a Rigid Body -3

- Angular acceleration

$${}^A\Omega_C = {}^A\Omega_B + {}^A_B R {}^B\Omega_C$$

↓ diff.

$$\begin{aligned} {}^A\dot{\Omega}_C &= {}^A\dot{\Omega}_B + \frac{d}{dt} {}^A_B R {}^B\Omega_C \\ &= {}^A\dot{\Omega}_B + {}^A_B R {}^B\dot{\Omega}_C + {}^A\Omega_B \times {}^A_B R {}^B\Omega_C \end{aligned}$$

Rigid Body Motion -1

□ Freshman Dynamics

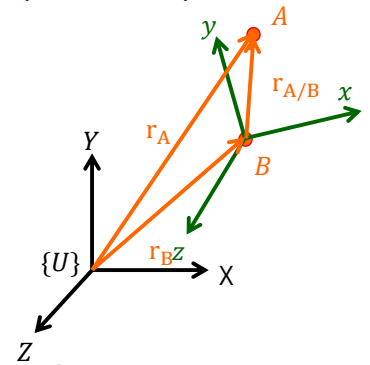
Following the materials described in Chap 5

$$\vec{v}_A = \vec{v}_B + \vec{v}_{rel} + \vec{\omega} \times \vec{r}_{A/B}$$

$$\vec{v}_A = (\dot{x}_B \hat{i} + \dot{y}_B \hat{j}) + (\dot{x}_{A/B} \hat{i} + \dot{y}_{A/B} \hat{j}) + \vec{\omega} \times (x_{A/B} \hat{i} + y_{A/B} \hat{j})$$

↓ diff.

$$\begin{aligned} \vec{a}_A = & (\ddot{x}_B \hat{i} + \ddot{y}_B \hat{j}) \\ & + (\ddot{x}_{A/B} \hat{i} + \ddot{y}_{A/B} \hat{j}) + \vec{\omega} \times (\dot{x}_{A/B} \hat{i} + \dot{y}_{A/B} \hat{j}) \\ & + \dot{\vec{\omega}} \times (x_{A/B} \hat{i} + y_{A/B} \hat{j}) \\ & + \vec{\omega} \times ((\dot{x}_{A/B} \hat{i} + \dot{y}_{A/B} \hat{j}) + \vec{\omega} \times (x_{A/B} \hat{i} + y_{A/B} \hat{j})) \end{aligned}$$



Rigid Body Motion -2

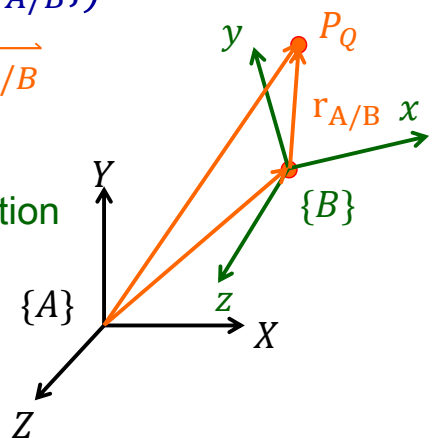
$$\begin{aligned} \vec{a}_A = & (\ddot{x}_B \hat{i} + \ddot{y}_B \hat{j}) \\ & + \dot{\vec{\omega}} \times (x_{A/B} \hat{i} + y_{A/B} \hat{j}) + \vec{\omega} \times (\vec{\omega} \times (x_{A/B} \hat{i} + y_{A/B} \hat{j})) \\ & + 2\vec{\omega} \times (\dot{x}_{A/B} \hat{i} + \dot{y}_{A/B} \hat{j}) + (\ddot{x}_{A/B} \hat{i} + \ddot{y}_{A/B} \hat{j}) \end{aligned}$$

$$\Rightarrow \vec{a}_A = \vec{a}_B + \underbrace{\dot{\vec{\omega}} \times \vec{r}_{A/B}}_{\text{Coriolis acceleration}} + \underbrace{\vec{\omega} \times \vec{\omega} \times \vec{r}_{A/B}}_{\text{"relative" acceleration}} + \underbrace{\ddot{x}_{A/B} \hat{i} + \ddot{y}_{A/B} \hat{j}}_{\text{relative acceleration}}$$

Coriolis acceleration "relative" acceleration

□ Thus

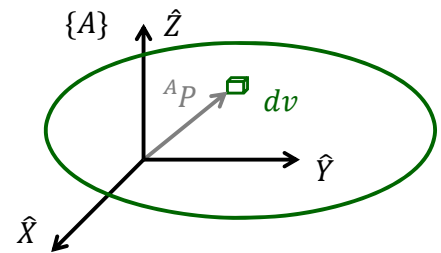
$$\begin{aligned} {}^A A_Q = & {}^A A_B \text{ ORG} \\ & + {}^A \dot{\Omega}_B \times {}^A R^B P_Q + {}^A \Omega_B \times ({}^A \Omega_B \times {}^A R^B P_Q) \\ & + 2 {}^A \Omega_B \times {}^A R^B V_Q + {}^A R^B A_Q \end{aligned}$$



Mass Distribution -1

- Inertia tensor relative to frame $\{A\}$

$${}^A I = \begin{bmatrix} I_{xx} & -I_{xy} & -I_{xz} \\ -I_{xy} & I_{yy} & -I_{yz} \\ -I_{xz} & -I_{yz} & I_{zz} \end{bmatrix}$$



Mass moment of inertia >0

$$I_{xx} = \iiint_V (y^2 + z^2) \rho dv$$

$$I_{yy} = \iiint_V (x^2 + z^2) \rho dv$$

$$I_{zz} = \iiint_V (x^2 + y^2) \rho dv$$

Mass product of inertia

$$I_{xy} = \iiint_V xy \rho dv$$

$$I_{xz} = \iiint_V xz \rho dv$$

$$I_{yz} = \iiint_V yz \rho dv$$

Mass Distribution -2

- Inertia tensor

- ◆ Constant real symmetric matrix (orthogonally diagonalizable)

its eigendecomposition (i. e., $M = V\Lambda V^{-1}$)

$${}^A I = \begin{bmatrix} I_{xx} & -I_{xy} & -I_{xz} \\ -I_{xy} & I_{yy} & -I_{yz} \\ -I_{xz} & -I_{yz} & I_{zz} \end{bmatrix} = R \begin{bmatrix} I_{XX} & 0 & 0 \\ 0 & I_{YY} & 0 \\ 0 & 0 & I_{ZZ} \end{bmatrix} R^T$$

↑
principal moment of inertia

Rotation matrix,
revealing directions of the principal axes

- ◆ $I_{xx} + I_{yy} + I_{zz} = \text{trace}({}^A I) = \text{constant}$
 - Trace is invariant under a similarity transformation
- ◆ If xy-plane is plane of symmetry, then $I_{xz} = I_{yz} = 0$

Mass Distribution -3

□ Parallel-axis Theorem

- ◆ Computing how the inertia tensor changes under translations of the reference coordinate system

$${}^A I_{zz} = {}^C I_{zz} + m(x_c^2 + y_c^2)$$

$${}^A I_{xy} = {}^C I_{xy} - mx_c y_c$$

C: at COM of the body

A: arbitrary frame

Vector-matrix form

$${}^A I = {}^C I + m[P_c^T P_c I_3 - P_c P_c^T]$$

$$P_c = [x_c \quad y_c \quad z_c]^T$$

COM relative to {A}

Newton's Equation and Euler's Equation

□ Newton's equation

$$F = \frac{d}{dt}(mv_c) = m\dot{v}_c$$

□ Euler's equation

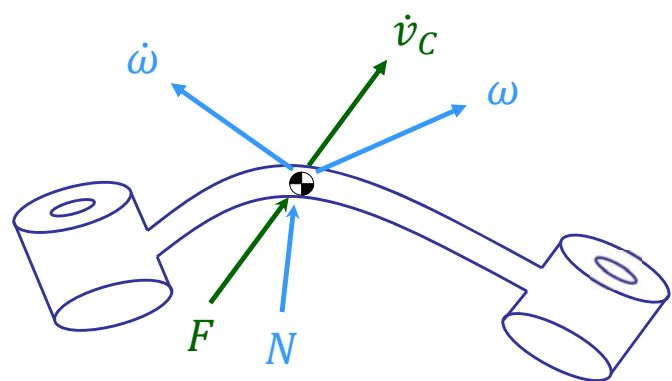
$$N = \frac{d}{dt}(I\omega)$$

Even if using inertial frame, it can change during motion

$$N = \underline{{}^C I} \dot{\omega} + \omega \times {}^C I \omega$$

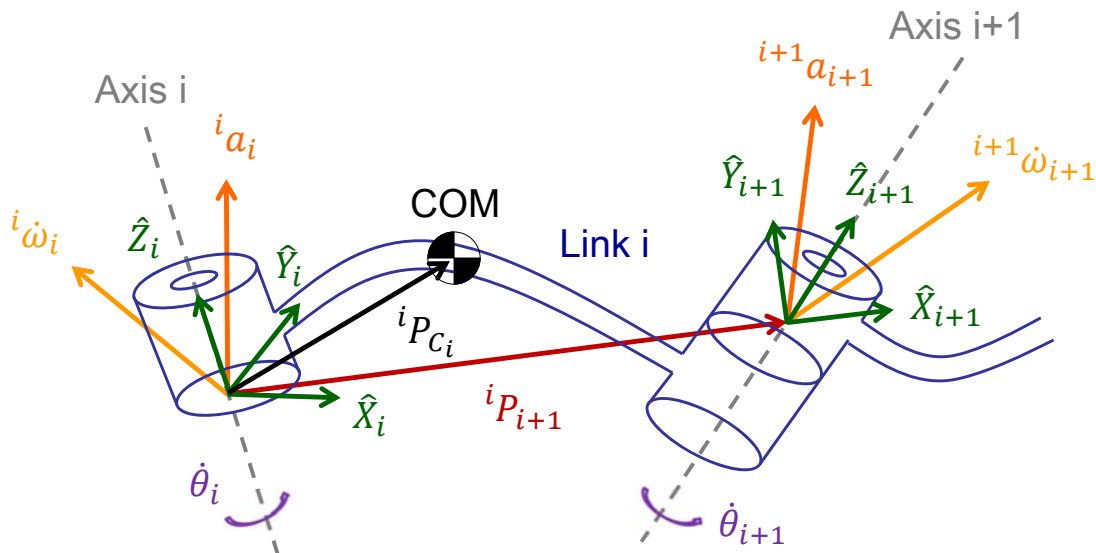
C: body frame, whose origin is located at COM

${}^C I$: constant matrix



Acceleration "Propagation" from Link to Link -1

- Strategy: Represent linear and angular accelerations of link i in frame $\{i\}$, and find their relationship to those of neighboring links



Acceleration "Propagation" from Link to Link -2

- Rotational Joint (Link $i+1$)
 - ◆ Angular acceleration propagation

In Chap 5

$${}^i\omega_{i+1} = {}^i\omega_i + \underbrace{{}_{i+1}{}^iR\dot{\theta}_{i+1}}_{} {}^{i+1}\hat{Z}_{i+1}$$

↓ diff.

$$\dot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1} = \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_{i+1} \end{bmatrix}$$

$${}^i\dot{\omega}_{i+1} = {}^i\dot{\omega}_i + {}^i\omega_i \times {}_{i+1}{}^iR\dot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1} + {}_{i+1}{}^iR\ddot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1}$$

↓ ${}^{i+1}{}_iR$

$${}^{i+i}\dot{\omega}_{i+1} = {}^{i+1}{}_iR {}^i\dot{\omega}_i + {}^{i+1}{}_iR {}^i\omega_i \times \dot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1} + \ddot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1}$$

Acceleration “Propagation” from Link to Link -3

- Linear acceleration propagation

$${}^i a_{i+1} = {}^i a_i + {}^i \dot{\omega}_i \times {}^i P_{i+1} + {}^i \omega_i \times ({}^i \omega_i \times {}^i P_{i+1})$$

$$\downarrow {}^{i+1}_i R$$

$${}^{i+1} a_{i+1} = {}^{i+1}_i R ({}^i a_i + {}^i \dot{\omega}_i \times {}^i P_{i+1} + {}^i \omega_i \times ({}^i \omega_i \times {}^i P_{i+1}))$$

Acceleration “Propagation” from Link to Link -4

- Prismatic joint (Link i+1)

- Angular acceleration propagation

$${}^i \dot{\omega}_i \xrightarrow{{}^{i+1}_i R} {}^{i+1} \dot{\omega}_{i+1} = {}^{i+1}_i R {}^i \dot{\omega}_i$$

- Linear acceleration propagation

$${}^i a_{i+1} = {}^i a_i + {}^i \dot{\omega}_i \times {}^i P_{i+1} + {}^i \omega_i \times ({}^i \omega_i \times {}^i P_{i+1}) \\ + 2 {}^i \omega_i \times {}^{i+1}_i R \dot{d}_{i+1} {}^{i+1} \hat{Z}_{i+1} + {}^{i+1}_i R \underline{\ddot{d}_{i+1}} {}^{i+1} \hat{Z}_{i+1}$$

$$\downarrow {}^{i+1}_i R$$

$$\ddot{d}_{i+1} {}^{i+1} \hat{Z}_{i+1} = \begin{bmatrix} 0 \\ 0 \\ \ddot{d}_{i+1} \end{bmatrix}$$

$${}^{i+1} a_{i+1} = {}^{i+1}_i R ({}^i a_i + {}^i \dot{\omega}_i \times {}^i P_{i+1} + {}^i \omega_i \times ({}^i \omega_i \times {}^i P_{i+1})) \\ + 2 {}^{i+1} \omega_{i+1} \times \dot{d}_{i+1} {}^{i+1} \hat{Z}_{i+1} + \ddot{d}_{i+1} {}^{i+1} \hat{Z}_{i+1}$$

Acceleration "Propagation" from Link to Link -5

□ COM

$${}^i a_{C_i} = {}^i a_i + {}^i \dot{\omega}_i \times {}^i P_{C_i} + {}^i \omega_i \times ({}^i \omega_i \times {}^i P_{C_i})$$

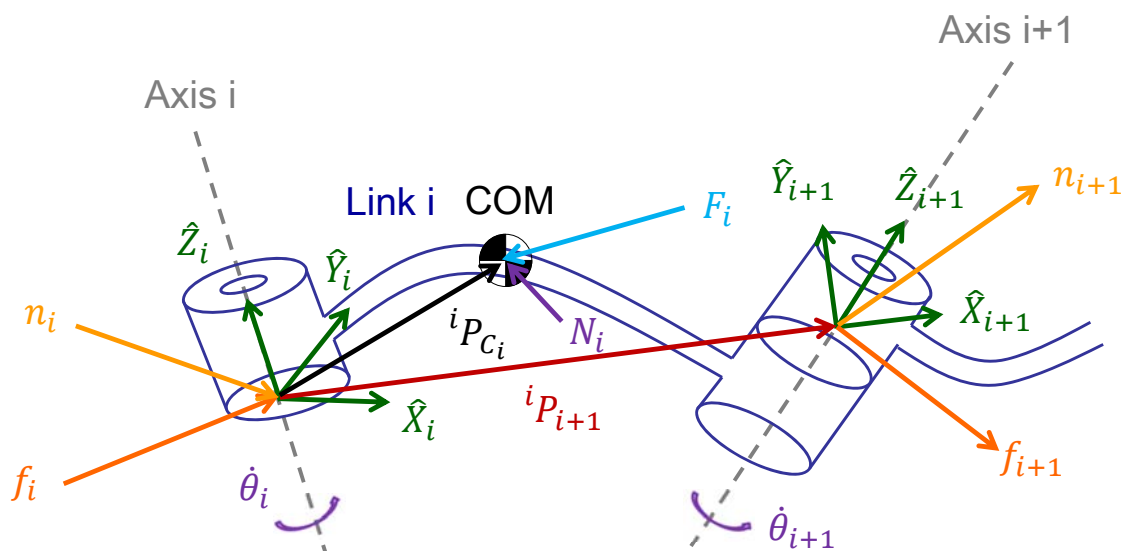
C_i : COM of the i^{th} link

Force Propagation from Link to Link -1

□ Inertia force and torque acting at the COM

$$F_i = m a_{C_i}$$

$$N_i = {}^c I \dot{\omega}_i + \omega_i \times {}^c I \omega_i$$

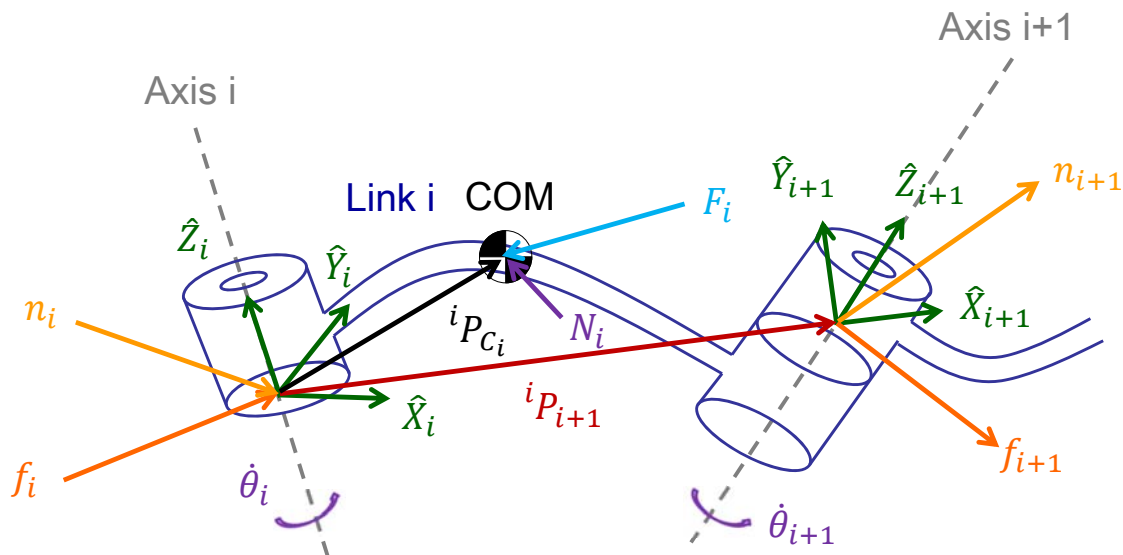


Force Propagation from Link to Link -2

□

$${}^i f_i = {}_{i+1}{}^i R^{i+1} f_{i+1} + {}^i F_i$$

$${}^i n_i = {}_{i+1}{}^i R^{i+1} n_{i+1} + {}^i N_i + {}^i P_{C_i} \times {}^i F_i + {}^i P_{i+1} \times {}_{i+1}{}^i R^{i+1} f_{i+1}$$



Force Propagation from Link to Link -3

□ Thus

- ◆ Revolute joint

$$\tau_i = {}^i n_i^T {}^i \hat{Z}_i$$

- ◆ Prismatic joint

$$\tau_i = {}^i f_i^T {}^i \hat{Z}_i$$

□ Comments

- ◆ Inclusion of gravity force: ${}^0 a_0 = g = 9.81 \text{ m/s}$
- ◆ A manipulator moving in free space: ${}^{N+1} f_{N+1} = 0$ ${}^{N+1} n_{N+1} = 0$

Iterative Newton-Euler Dynamic Formulation

- Outward iterations
 - ◆ Link 1 to link n
 - ◆ Velocities and accelerations
- Inward iterations
 - ◆ Link n to link 1
 - ◆ Forces and torques
- Revolute joint vs. prismatic joint: Choose correct equations
- General structure, can be applied to any manipulator
- Easy for numerical computation

Example: A RR Manipulator -1

- Conditions:

$${}^1P_{C_1} = l_1 \hat{X}_1 \quad c_1 I_1 = 0$$

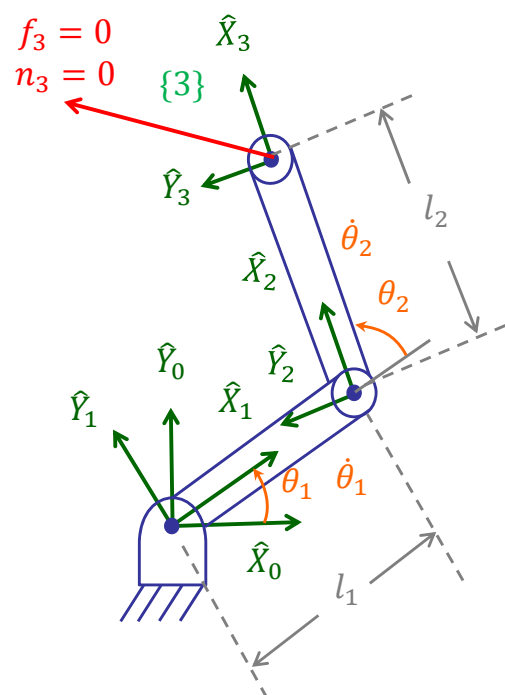
$${}^2P_{C_2} = l_2 \hat{X}_2 \quad c_2 I_2 = 0$$

$$m_1, m_2$$

$$\omega_0 = 0 \quad {}^0v_0 = g \hat{Y}_0$$

$$\dot{\omega}_0 = 0$$

$${}_{i+1}^i R = \begin{bmatrix} c_{i+1} & -s_{i+1} & 0 \\ s_{i+1} & c_{i+1} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$



Example: A RR Manipulator -2

- Velocity and acceleration propagations

$${}^1\omega_1 = {}^0R^0\omega_0 + \dot{\theta}_1 {}^1\hat{Z}_1 = \dot{\theta}_1 {}^1\hat{Z}_1 = \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_1 \end{bmatrix}$$

$${}^1\dot{\omega}_1 = {}^0R^0\dot{\omega}_0 + {}^0R^0\omega_0 \times \dot{\theta}_1 {}^1\hat{Z}_1 + \ddot{\theta}_1 \hat{Z}_1 = \begin{bmatrix} 0 \\ 0 \\ \ddot{\theta}_1 \end{bmatrix}$$

$${}^1a_1 = {}^0R({}^0a_0 + {}^0\omega_0 \times {}^0P_1 + {}^0\omega_0 \times ({}^0\omega_0 \times {}^0P_1)) = \begin{bmatrix} c_1 & s_1 & 0 \\ -s_1 & c_1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ g \\ 0 \end{bmatrix}$$

$$\begin{aligned} {}^1a_{c_1} &= {}^1a_1 + {}^1\dot{\omega}_1 \times {}^1P_{c_1} + {}^1\omega_1 \times ({}^1\omega_1 \times {}^1P_{c_1}) \\ &= \begin{bmatrix} gs_1 \\ gc_1 \\ 0 \end{bmatrix} + \begin{bmatrix} 0 \\ l_1\ddot{\theta}_1 \\ 0 \end{bmatrix} + \begin{bmatrix} -l_1\dot{\theta}_1^2 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} -l_1\dot{\theta}_1^2 + gs_1 \\ l_1\ddot{\theta}_1 + gc_1 \\ 0 \end{bmatrix} \end{aligned}$$

Example: A RR Manipulator -3

$${}^1F_1 = m {}^1a_{c_1} \begin{bmatrix} -m_1 l_1 \dot{\theta}_1^2 + m_1 g s_1 \\ m_1 l_1 \ddot{\theta}_1 + m_1 g c_1 \\ 0 \end{bmatrix}$$

$${}^1N_1 = {}^{c_1}I {}^1\dot{\omega}_1 + {}^1\omega_1 \times {}^{c_1}I {}^1\omega_1 = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$${}^2\omega_2 = {}^2R {}^1\omega_1 + \dot{\theta}_2 {}^2\hat{Z}_2 = \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_1 + \dot{\theta}_2 \end{bmatrix}$$

$${}^2\dot{\omega}_2 = {}^2R {}^1\dot{\omega}_1 + {}^2R {}^1\omega_1 \times \dot{\theta}_2 {}^2\hat{Z}_2 + \ddot{\theta}_2 {}^2\hat{Z}_2 = \begin{bmatrix} 0 \\ 0 \\ \ddot{\theta}_1 + \ddot{\theta}_2 \end{bmatrix}$$

$$\begin{aligned} {}^2a_2 &= {}^2R({}^1a_1 + {}^1\dot{\omega}_1 \times {}^1P_2 + {}^1\omega_1 \times ({}^1\omega_1 \times {}^1P_2)) = \\ &= \begin{bmatrix} c_2 & s_2 & 0 \\ -s_2 & c_2 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} -l_1\dot{\theta}_1^2 + gs_1 \\ l_1\ddot{\theta}_1 + gc_1 \\ 0 \end{bmatrix} = \begin{bmatrix} l_1\ddot{\theta}_1 s_2 - l_1\dot{\theta}_1^2 c_2 + gs_{12} \\ l_1\ddot{\theta}_1 c_2 + l_1\dot{\theta}_1^2 s_2 + gc_{12} \\ 0 \end{bmatrix} \end{aligned}$$

Example: A RR Manipulator -4

$$\begin{aligned}
 {}^2a_{C_2} &= {}^2a_2 + {}^2\dot{\omega}_2 \times {}^2P_{C_2} + {}^2\omega_2 \times ({}^2\omega_2 \times {}^2P_{C_2}) \\
 &= \begin{bmatrix} l_1\ddot{\theta}_1 s_2 - l_1\dot{\theta}_1^2 c_2 + g s_{12} \\ l_1\ddot{\theta}_1 c_2 + l_1\dot{\theta}_1^2 s_2 + g c_{12} \\ 0 \end{bmatrix} + \begin{bmatrix} 0 \\ l_2(\ddot{\theta}_1 + \ddot{\theta}_2) \\ 0 \end{bmatrix} + \begin{bmatrix} -l_2(\dot{\theta}_1 + \dot{\theta}_2)^2 \\ 0 \\ 0 \end{bmatrix} \\
 {}^2F_2 &= m {}^2a_{C_2} = \begin{bmatrix} m_2 l_1 \ddot{\theta}_1 s_2 - m_2 l_1 \dot{\theta}_1^2 c_2 + m_2 g s_{12} - m_2 l_2 (\dot{\theta}_1 + \dot{\theta}_2)^2 \\ m_2 l_1 \ddot{\theta}_1 c_2 + m_2 l_1 \dot{\theta}_1^2 s_2 + m_2 g c_{12} + m_2 l_2 (\ddot{\theta}_1 + \ddot{\theta}_2) \\ 0 \end{bmatrix} \\
 {}^2N_2 &= {}^{C_2}I {}^2\dot{\omega}_2 + {}^2\omega_2 \times {}^{C_2}I {}^2\omega_2 = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}
 \end{aligned}$$

Example: A RR Manipulator -5

□ Force and torque propagations

$${}^2f_2 = {}^2_3R {}^3f_3 + {}^2F_2 = {}^2F_2$$

$$\begin{aligned}
 {}^2n_2 &= {}^2_3R {}^3n_3 + {}^2N_2 + {}^2P_{C_2} \times {}^2F_2 + {}^2P_3 \times {}^2_3R {}^3f_3 \\
 &= \begin{bmatrix} 0 \\ 0 \\ m_2 l_1 l_2 c_2 \ddot{\theta}_1 + m_2 l_1 l_2 s_2 \dot{\theta}_1^2 + m_2 l_2 g c_{12} + m_2 l_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2) \end{bmatrix}
 \end{aligned}$$

$$\begin{aligned}
 {}^1f_1 &= {}^1_2R {}^2f_2 + {}^1F_1 \\
 &= \begin{bmatrix} c_2 & -s_2 & 0 \\ s_2 & c_2 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} m_2 l_1 s_2 \ddot{\theta}_1 - m_2 l_1 c_2 \dot{\theta}_1^2 + m_2 g s_{12} - m_2 l_2 (\dot{\theta}_1 + \dot{\theta}_2)^2 \\ m_2 l_1 c_2 \ddot{\theta}_1 + m_2 l_1 s_2 \dot{\theta}_1^2 + m_2 g c_{12} + m_2 l_2 (\ddot{\theta}_1 + \ddot{\theta}_2) \\ 0 \end{bmatrix} \\
 &\quad + \begin{bmatrix} -m_1 l_1 \dot{\theta}_1^2 + m_1 g s_1 \\ m_1 l_1 \ddot{\theta}_1 + m_1 g c_1 \\ 0 \end{bmatrix}
 \end{aligned}$$

Example: A RR Manipulator -6

$$\begin{aligned}
 {}^1n_1 &= {}^1_2R {}^2n_2 + {}^1N_1 + {}^1P_{C_1} \times {}^1F_1 + {}^1P_2 \times {}^1_2R {}^2f_2 \\
 &= \begin{bmatrix} 0 \\ 0 \\ m_2l_1l_2c_2\ddot{\theta}_1 + m_2l_1l_2s_2\dot{\theta}_1^2 + m_2l_2gc_{12} + m_2l_2^2(\ddot{\theta}_1 + \ddot{\theta}_2) \end{bmatrix} \\
 &\quad + \begin{bmatrix} 0 \\ 0 \\ m_1l_1^2\ddot{\theta}_1 + m_1l_1gc_1 \end{bmatrix} \\
 &\quad + \begin{bmatrix} 0 \\ 0 \\ m_2l_1^2\ddot{\theta}_1 - m_2l_1l_2s_2(\dot{\theta}_1 + \dot{\theta}_2)^2 + m_2l_1gs_2s_{12} + m_2l_1l_2c_2(\ddot{\theta}_1 + \ddot{\theta}_2) + m_2l_1gc_2c_{12} \end{bmatrix}
 \end{aligned}$$

Example: A RR Manipulator -7

□ Joint torques

$$\begin{aligned}
 \tau_1 &= {}^1n_1^T {}^1\widehat{Z}_1 \\
 &= m_2l_2^2(\ddot{\theta}_1 + \ddot{\theta}_2) + m_2l_1l_2c_2(2\dot{\theta}_1 + \dot{\theta}_2) + (m_1 + m_2)l_1^2\ddot{\theta}_1 \\
 &\quad - m_2l_1l_2s_2\dot{\theta}_2^2 - 2m_2l_1l_2s_2\dot{\theta}_1\dot{\theta}_2 + m_2l_2gc_{12} + (m_1 + m_2)l_1gc_1
 \end{aligned}$$

$$\begin{aligned}
 \tau_2 &= {}^2n_2^T {}^2\widehat{Z}_2 \\
 &= m_2l_1l_2c_2\ddot{\theta}_1 + m_2l_1l_2s_2\dot{\theta}_1^2 + m_2l_2gc_{12} + m_2l_2^2(\ddot{\theta}_1 + \ddot{\theta}_2)
 \end{aligned}$$

The Structure of Dynamic Equations -1

- The state-space equation

$$\tau = M(\Theta)\ddot{\Theta} + V(\Theta, \dot{\Theta}) + G(\Theta)$$

$n \times 1$ $n \times n$ $n \times 1$ $n \times 1$ $n \times 1$
 Mass Centrifugal gravity
 matrix Coriolis

- Revisit the RR manipulator

$$M(\Theta) = \begin{bmatrix} l_2^2 m_2 + 2l_1 l_2 m_2 c_2 + l_1^2 (m_1 + m_2) & l_2^2 m_2 + l_1 l_2 m_2 c_2 \\ l_2^2 m_2 + l_1 l_2 m_2 c_2 & l_2^2 m_2 \end{bmatrix}$$

$$V(\Theta, \dot{\Theta}) = \begin{bmatrix} -m_2 l_1 l_2 s_2 \dot{\theta}_2^2 - 2m_2 l_1 l_2 s_2 \dot{\theta}_1 \dot{\theta}_2 \\ m_2 l_1 l_2 s_2 \dot{\theta}_1^2 \end{bmatrix}$$

$$G(\Theta) = \begin{bmatrix} m_2 l_2 g c_{12} + (m_1 + m_2) l_1 g c_1 \\ m_2 l_2 g c_{12} \end{bmatrix}$$

The Structure of Dynamic Equations -2

- The configuration-space equation

$$\tau = M(\Theta)\ddot{\Theta} + B(\Theta)[\dot{\Theta}\dot{\Theta}] + C(\Theta)[\dot{\Theta}^2] + G(\Theta)$$

$n \times 1$ $n \times n$ $n \times \frac{n(n-1)}{2}$ $n \times n$ $n \times 1$
 Mass Coriolis Centrifugal gravity
 matrix
 $n \times 1$

$$[\dot{\Theta}\dot{\Theta}] = [\dot{\theta}_1 \dot{\theta}_2 \quad \dot{\theta}_1 \dot{\theta}_3 \quad \dots \quad \dot{\theta}_{n-1} \dot{\theta}_n]^T$$

$\frac{n(n-1)}{2} \times 1$

$$[\dot{\Theta}^2] = [\dot{\theta}_1^2 \quad \dot{\theta}_2^2 \quad \dots \quad \dot{\theta}_n^2]^T$$

$n \times 1$

The Structure of Dynamic Equations -3

- Revisit the RR manipulator

$$V(\Theta, \dot{\Theta}) = \begin{bmatrix} -m_2 l_1 l_2 s_2 \dot{\theta}_2^2 - 2m_2 l_1 l_2 s_2 \dot{\theta}_1 \dot{\theta}_2 \\ m_2 l_1 l_2 s_2 \dot{\theta}_1^2 \end{bmatrix} = B(\Theta) [\dot{\Theta} \dot{\Theta}] + C(\Theta) [\dot{\Theta}^2]$$

$$B(\Theta) = \begin{bmatrix} -2m_2 l_1 l_2 s_2 \\ 0 \end{bmatrix} \quad [\dot{\Theta} \dot{\Theta}] = [\dot{\theta}_1 \dot{\theta}_2]$$

$$C(\Theta) = \begin{bmatrix} 0 & -m_2 l_1 l_2 s_2 \\ m_2 l_1 l_2 s_2 & 0 \end{bmatrix} \quad [\dot{\Theta}^2] = \begin{bmatrix} \dot{\theta}_1^2 \\ \dot{\theta}_2^2 \end{bmatrix}$$

Lagrangian Formulation of Manipulator Dynamics -1

- Newton-Euler: Force-moment-based analysis

Lagrange: Energy-based analysis

- Of course, for a system, both methods should yield the same equations of motion

Lagrangian Formulation of Manipulator Dynamics -2

- Kinetic energy

$$k_i = \frac{1}{2} m_i v_{c_i}^T v_{c_i} + \frac{1}{2} {}^i \omega_i^T C_i I_i {}^i \omega_i$$

$$k = \sum_{i=1}^n k_i \quad k = k(\Theta, \dot{\Theta}) = \frac{1}{2} \dot{\Theta}^T M(\Theta) \dot{\Theta}$$

- Potential energy

$$u_i = -m_i {}^0 g^T {}^0 P_{C_i} + \underline{u_{ref_i}}$$

Shift the zero reference height

$$u = \sum_{i=1}^n u_i \quad u = u(\Theta)$$

Lagrangian Formulation of Manipulator Dynamics -3

- Lagrangian

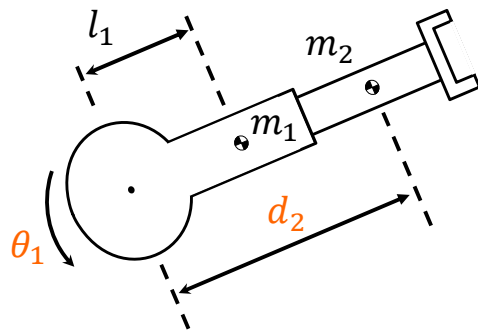
$$\mathcal{L}(\Theta, \dot{\Theta}) = k(\Theta, \dot{\Theta}) - u(\Theta)$$

- Equation of motion for the manipulator

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\Theta}} - \frac{\partial \mathcal{L}}{\partial \Theta} = \tau$$

$$\frac{d}{dt} \frac{\partial k}{\partial \dot{\Theta}} - \frac{\partial k}{\partial \Theta} + \frac{\partial u}{\partial \Theta} = \tau$$

Example: An RP Manipulator -1



$$c_1 I_1 = \begin{bmatrix} I_{xx1} & 0 & 0 \\ 0 & I_{yy1} & 0 \\ 0 & 0 & I_{zz1} \end{bmatrix}$$

$$c_2 I_2 = \begin{bmatrix} I_{xx2} & 0 & 0 \\ 0 & I_{yy2} & 0 \\ 0 & 0 & I_{zz2} \end{bmatrix}$$

◆ Kinetic energy

$$k_1 = \frac{1}{2} m_1 l_1^2 \dot{\theta}_1^2 + \frac{1}{2} I_{zz1} \dot{\theta}_1^2$$

$$k_2 = \frac{1}{2} m_2 (d_2^2 \dot{\theta}_1^2 + \dot{d}_2^2) + \frac{1}{2} I_{zz2} \dot{\theta}_1^2$$

$$k(\theta, \dot{\theta}) = \frac{1}{2} (m_1 l_1^2 + I_{zz1} + I_{zz2} + m_2 d_2^2) \dot{\theta}_1^2 + \frac{1}{2} m_2 \dot{d}_2^2$$

Example: An RP Manipulator -2

◆ Potential energy

$$u_1 = m_1 g l_1 \sin \theta_1 + m_1 g l_1$$

$$u_2 = m_2 g d_2 \sin \theta_1 + m_2 g d_{2max}$$

$$u(\Theta) = (m_1 l_1 + m_2 d_2) g \sin \theta_1 + \underline{m_1 g l_1 + m_2 g d_{2max}}$$

Shift the zero reference height

◆ Lagrangian

$$\frac{\partial k}{\partial \dot{\Theta}} = \begin{bmatrix} (m_1 l_1^2 + I_{zz1} + I_{zz2} + m_2 d_2^2) \dot{\theta}_1 \\ m_2 \dot{d}_2 \end{bmatrix}$$

$$\frac{\partial k}{\partial \Theta} = \begin{bmatrix} 0 \\ m_2 d_2 \dot{\theta}_1^2 \end{bmatrix}$$

$$\frac{\partial u}{\partial \Theta} = \begin{bmatrix} (m_1 l_1 + m_2 d_2) g \cos \theta_1 \\ m_2 g \sin \theta_1 \end{bmatrix}$$

Example: An RP Manipulator -3

- ◆ Equations of motion

$$\tau_1 = (m_1 l_1^2 + I_{zz1} + I_{zz2} + m_2 d_2^2) \ddot{\theta}_1 + 2m_2 d_2 \dot{\theta}_1 \dot{d}_2 + (m_1 l_1 + m_2 d_2) g \cos \theta_1$$

$$\tau_2 = m_2 \ddot{d}_2 - m_2 d_2 \dot{\theta}_1^2 + m_2 g \sin \theta_1$$

state-space representation

$$\tau = M(\theta) \ddot{\theta} + V(\theta, \dot{\theta}) + G(\theta)$$

$$M(\theta) = \begin{bmatrix} m_1 l_1^2 + I_{zz1} + I_{zz2} + m_2 d_2^2 & 0 \\ 0 & m_2 \end{bmatrix}$$

$$V(\theta, \dot{\theta}) = \begin{bmatrix} 2m_2 d_2 \dot{\theta}_1 \dot{d}_2 \\ -m_2 d_2 \dot{\theta}_1^2 \end{bmatrix}$$

$$G(\theta) = \begin{bmatrix} (m_1 l_1 + m_2 d_2) g \cos \theta_1 \\ m_2 g \sin \theta_1 \end{bmatrix}$$

Manipulator Dynamics in Cartesian Space -1

- Dynamic equations

- ◆ In joint space

$$\tau = M(\theta) \ddot{\theta} + V(\theta, \dot{\theta}) + G(\theta)$$

- ◆ In Cartesian space

$$F = M_x(\theta) \ddot{X} + V_x(\theta, \dot{\theta}) + G_x(\theta)$$

- Formulation

$$\tau = J^T(\theta) F$$

$$F = J^{-T} \tau = J^{-T} M(\theta) \ddot{\theta} + J^{-T} V(\theta, \dot{\theta}) + J^{-T} G(\theta)$$

$$\dot{X} = J \dot{\theta} \quad \ddot{X} = j \dot{\theta} + J \ddot{\theta} \quad \ddot{\theta} = J^{-1} \ddot{X} - J^{-1} j \dot{\theta}$$

Manipulator Dynamics in Cartesian Space -2

$$F = J^{-T} M(\Theta) J^{-1} \ddot{X} - J^{-T} M(\Theta) J^{-1} \dot{J} \dot{\Theta} + J^{-T} V(\Theta, \dot{\Theta}) + J^{-T} G(\Theta) \\ = M_x(\Theta) \ddot{X} + V_x(\Theta, \dot{\Theta}) + G_x(\Theta)$$

$$M_x(\Theta) = J^{-T}(\Theta) M(\Theta) J^{-1}(\Theta)$$

$$V_x(\Theta, \dot{\Theta}) = J^{-T}(\Theta) (V(\Theta, \dot{\Theta}) - M(\Theta) J^{-1}(\Theta) \dot{J}(\Theta) \dot{\Theta})$$

$$G_x(\Theta) = J^{-T}(\Theta) G(\Theta)$$

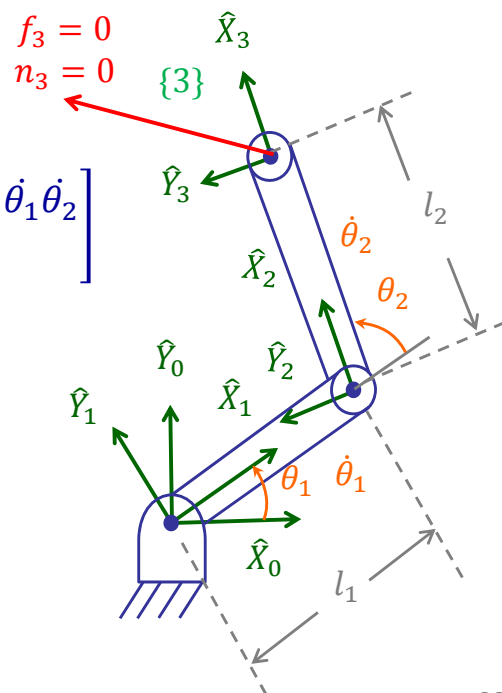
Revisit Example: A RR Manipulator -1

- In joint space

$$M(\Theta) = \begin{bmatrix} l_2^2 m_2 + 2l_1 l_2 m_2 c_2 + l_1^2 (m_1 + m_2) & l_2^2 m_2 + l_1 l_2 m_2 c_2 \\ l_2^2 m_2 + l_1 l_2 m_2 c_2 & l_2^2 m_2 \end{bmatrix}$$

$$V(\Theta, \dot{\Theta}) = \begin{bmatrix} -m_2 l_1 l_2 s_2 \dot{\theta}_2^2 - 2m_2 l_1 l_2 s_2 \dot{\theta}_1 \dot{\theta}_2 \\ m_2 l_1 l_2 s_2 \dot{\theta}_1^2 \end{bmatrix}$$

$$G(\Theta) = \begin{bmatrix} m_2 l_2 g c_{12} + (m_1 + m_2) l_1 g c_1 \\ m_2 l_2 g c_{12} \end{bmatrix}$$



Revisit Example: A RR Manipulator -2

□ Jacobian

$$J(\Theta) = \begin{bmatrix} l_1 s_2 & 0 \\ l_1 c_2 + l_2 & l_2 \end{bmatrix} \quad J^{-1}(\Theta) = \frac{1}{l_1 l_2 s_2} \begin{bmatrix} l_2 & 0 \\ -l_1 c_2 - l_2 & l_1 s_2 \end{bmatrix}$$

$$j(\Theta) = \begin{bmatrix} l_1 c_2 \dot{\theta}_2 & 0 \\ -l_1 s_2 \dot{\theta}_2 & 0 \end{bmatrix}$$

□ In Cartesian space

$$M_x(\Theta) = J^{-T}(\Theta)M(\Theta)J^{-1}(\Theta) = \begin{bmatrix} m_2 + \frac{m_1}{s_2^2} & 0 \\ 0 & m_2 \end{bmatrix}$$

$$V_x(\Theta, \dot{\Theta}) = J^{-T}(\Theta)(V(\Theta, \dot{\Theta}) - M(\Theta)J^{-1}(\Theta)j(\Theta)\dot{\Theta})$$
$$= \begin{bmatrix} -(m_2 l_1 c_2 + m_2 l_2) \dot{\theta}_1^2 - m_2 l_2 \dot{\theta}_2^2 - (2m_2 l_2 + m_2 l_1 c_2 + m_1 l_1 \frac{c_2}{s_2^2}) \dot{\theta}_1 \dot{\theta}_2 \\ m_2 l_1 s_2 \dot{\theta}_1^2 + m_2 l_1 s_2 \dot{\theta}_1 \dot{\theta}_2 \end{bmatrix}$$

$$G_x(\Theta) = J^{-T}(\Theta)G(\Theta) = \begin{bmatrix} m_1 g \frac{c_1}{s_2} + m_2 g s_{12} \\ m_2 g c_{12} \end{bmatrix}$$

Torque Equation

□ In Cartesian space

$$\tau = J^T(\Theta)F = J^T(\Theta)(M_x(\Theta)\ddot{X} + V_x(\Theta, \dot{\Theta}) + G_x(\Theta))$$

$$\tau = J^T(\Theta)M_x(\Theta)\ddot{X} + B_x(\Theta)[\dot{\Theta}\dot{\Theta}] + C_x(\Theta)[\dot{\Theta}^2] + G(\Theta)$$

$$\Rightarrow J^T(\Theta)V_x(\Theta, \dot{\Theta}) = B_x(\Theta)[\dot{\Theta}\dot{\Theta}] + C_x(\Theta)[\dot{\Theta}^2]$$

Revisit Example: A RR Manipulator

$$J^T(\Theta)V_x(\Theta, \dot{\Theta}) = B_x(\Theta)[\dot{\Theta}\dot{\Theta}] + C_x(\Theta)[\dot{\Theta}^2]$$
$$= \begin{bmatrix} l_1 s_2 & l_1 c_2 + l_2 \\ 0 & l_2 \end{bmatrix} \begin{bmatrix} -(m_2 l_1 c_2 + m_2 l_2) \dot{\theta}_1^2 - m_2 l_2 \dot{\theta}_2^2 - (2m_2 l_2 + m_2 l_1 c_2 + m_1 l_1 \frac{c_2}{s_2^2}) \dot{\theta}_1 \dot{\theta}_2 \\ m_2 l_1 s_2 \dot{\theta}_1^2 + l_1 m_2 s_2 \dot{\theta}_1 \dot{\theta}_2 \end{bmatrix}$$

$$B_x(\Theta) = \begin{bmatrix} -m_1 l_1^2 \frac{c_2}{s_2} - m_2 l_1 l_2 s_2 \\ m_2 l_1 l_2 s_2 \end{bmatrix}$$

$$C_x(\Theta) = \begin{bmatrix} 0 & -m_2 l_1 l_2 s_2 \\ m_2 l_1 l_2 s_2 & 0 \end{bmatrix}$$

Friction

- Viscous friction

$$\tau_{friction} = c\dot{\theta}$$

- Coulomb friction

$$\tau_{friction} = c \operatorname{sgn}\dot{\theta}$$

$$\dot{\theta} = 0, c = \text{“static coefficient”}$$

$$\dot{\theta} \neq 0, c = \text{“dynamic coefficient”}$$

- Questions?

